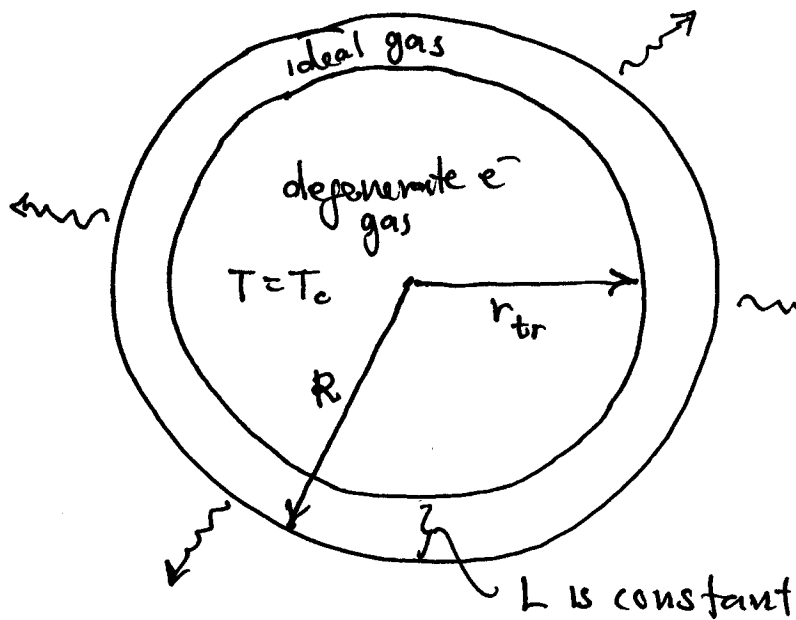


WD cooling

assumptions:

- core is degenerate, contains nearly all M , R
- thin radiative envelope surrounds the core (ideal gas)
- transition b/w core and envelope is abrupt, at $r = r_{tr}$



- conduction dominates in core — nearly isothermal
- $T_{tr} = T_{core}$
- no ν losses, ϵ_{nuc} , or contraction

recall the radiative envelope discussion when we covered polytropes

$$\text{completely radiative: } \nabla = \nabla_{rad} = \frac{d \log T}{d \log P} = \frac{3}{16\pi a c G} \frac{P \bar{\kappa}}{T^4} \frac{L_*}{M_*}$$

we took: P as ideal gas

$$\kappa = \kappa_0 \rho^0 T^{-s} = \kappa_g P^0 T^{-\nu-s}$$

$$\kappa_g = \kappa_0 \left(\frac{\mu m_u}{k} \right)^0$$

thin envelope

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We rewrite this as

$$T^4 d \log T = \left(\frac{3 \bar{k}}{16 \pi a c G} \right) P d \log P$$

or

$$T^{\nu+s+3} dT = \left(\frac{3 k_g L_*}{16 \pi a c G M_*} \right) P^\nu dP$$

and we integrated from T_0, P_0 to $T(r) \geq T_0, P(r) \geq P_0$
(inward)

giving

$$P^{\nu+1} \left[1 - \left(\frac{P_0}{P} \right)^{\nu+1} \right] = \frac{16 \pi a c G}{3 k_g} \frac{M_*}{L_*} \frac{\nu+1}{\nu+s+4} T^{\nu+s+4} \left[1 - \left(\frac{T_0}{T} \right)^{\nu+s+4} \right]$$

w/ $\nu+s+4 \neq 0$

Now, if $\nu+s+4 > 0$ and $\nu+1 > 0$, then at a depth beneath the surface, $P(r) \gg P_0, T(r) \gg T_0$, giving

$$P^{\nu+1} = \frac{16 \pi a c G}{3 k_g} \frac{M_*}{L_*} \frac{\nu+1}{\nu+s+4} T^{\nu+s+4}$$

this has the form

$$P = K' T^{1+n_{\text{eff}}} \quad \text{w/} \quad n_{\text{eff}} = \frac{s+3}{\nu+1}$$

and

$$K' = \left[\frac{1}{1+n_{\text{eff}}} \frac{16 \pi a c G M_*}{3 k_g L_*} \left(\frac{k}{\mu m_u} \right)^\nu \right]^{\frac{1}{\nu+1}}$$

k_0 not k_g

If we have a constant composition then $K' = \text{constant}$

and

$$p = \frac{\rho kT}{\mu m_0} = K' T^{1+n_{\text{eff}}} \rightarrow \rho \sim T^{-n_{\text{eff}}}$$

$$\therefore T \sim \rho^{-1/n_{\text{eff}}} \quad \text{and} \quad P \sim \rho T \sim \rho^{1+1/n_{\text{eff}}} \quad (\text{a polytrope!})$$

4.

What is the transition from core to envelope?

$$E_f \sim kT$$

$$\frac{p_f^2}{2m_e} \sim kT$$

or using $x_F = \frac{p_F}{m_e c}$

$$x_F^2 \sim \frac{2kT}{m_e c^2}$$

from our number density constraint, we know

$$n_e = \frac{8\pi}{3} \left(\frac{m_e c}{h} \right)^3 x_F^3 \quad (\text{HKT 3.50})$$

$$\therefore n_e \sim \left(\frac{2kT}{m_e c^2} \right)^{3/2} \frac{8\pi}{3} \left(\frac{m_e c}{h} \right)^3$$

$$\frac{\rho}{\mu_e m_u} \sim \left(\frac{2kT}{m_e c^2} \right)^{3/2} \frac{8\pi}{3} \left(\frac{m_e c}{h} \right)^3$$

$$\left(\frac{\rho}{\mu_e m_u} \right)^{2/3} \sim \frac{2kT}{m_e c^2} \left(\frac{8\pi}{3} \right)^{2/3} \left(\frac{m_e c}{h} \right)^2$$

or $kT \sim \left(\frac{\rho}{\mu_e} \right)^{2/3} \frac{1}{2m_e} \left(\frac{3h^3}{8\pi} \right)^{2/3} \frac{1}{m_u^{2/3}}$

this is $T \sim 2.9 \times 10^5 \left(\frac{\rho}{\mu_e} \right)^{2/3} \quad [\text{CGS}]$

or $\left(\frac{\rho}{\mu_e} \right) \sim \underbrace{6 \times 10^{-9}}_{\equiv \alpha} T^{3/2} \quad [\text{CGS}]$

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The transition occurs when

$$T_{tr} \sim T_c \quad (\text{since conduction is so efficient})$$

requiring that the pressure is continuous across the transition,

$$p_{tr} = K' T_{tr}^{1+n_{eff}} = \frac{p_{tr} k T_{tr}}{\mu m_u} = \alpha \mu_e \frac{k}{\mu m_u} T_{tr}^{1+3/2}$$

↑
use $p_{tr} \sim \alpha \mu_e T_{tr}^{3/2}$

We'll assume that the envelope composition is the same as the core (probably not true)

K' depends on L_* , M_* , μ , and κ

Bound-free opacity gives

$$\kappa_{bf} \sim 4 \times 10^{25} \rho T^{-7/2} \text{ cm}^2 \text{ g}^{-1}$$

$$\text{so } \nu = 1, \quad s = +7/2, \quad \text{and } n_{eff} = \frac{s+3}{\nu+1} = 3.25$$

$$K' = \left[\frac{1}{1+n_{eff}} \frac{16\pi a c G M_*}{3 k_0 L_*} \left(\frac{k}{\mu m_u} \right) \right]^{1/2}$$
$$= 8 \times 10^{-15} \mu^{-1/2} \left(\frac{M_*}{M_\odot} \right)^{1/2} \left(\frac{L_*}{L_\odot} \right)^{-1/2} \quad [\text{CGS}]$$

taking $\mu_e = 2$, we have

$$K' T_{tr}^{1+3.25} = \alpha (2) \frac{k}{\mu m_u} T_{tr}^{1+3/2}$$

$$8 \times 10^{-15} \mu^{-1/2} \left(\frac{M_*}{M_\odot} \right)^{1/2} \left(\frac{L_*}{L_\odot} \right)^{-1/2} T_{tr}^{17/4} = 6 \times 10^{-9} \cdot 2 \cdot \frac{k}{\mu m_u} T_{tr}^{5/2}$$

6. Together this gives

$$\frac{L_{\star}}{L_{\odot}} = 6.5 \times 10^{-29} \mu \left(\frac{M_{\star}}{M_{\odot}} \right) T_{tr}^{7/2} \quad [CGS]$$

Note: we can get the thickness of the envelope using the polytrope structure we derived earlier

Assume no internal heat sources, then $R \sim \text{const}$

most of the heat is in the ions:

$$c_v = \left. \frac{\partial e}{\partial T} \right|_p = \frac{3}{2} \frac{k}{\mu_{\pm} m_p} \sim \frac{1.2 \times 10^8}{\mu_{\pm}} \text{ erg/g/K}$$

the luminosity is just

$$L_{\star} = - \frac{dE_{ion}}{dt} = - \underbrace{c_v M_{\star}}_{\substack{\text{core has all the} \\ \text{mass and const } T}} \frac{dT_{tr}}{dt}$$

equating, starting w/ previous result

$$L_{\star} = 6.5 \times 10^{-29} \mu \left(\frac{M_{\star}}{M_{\odot}} \right) L_{\odot} T_{tr}^{7/2}$$

$$\frac{dL_{\star}}{dt} = \frac{7}{2} \eta \mu M_{\star} T_{tr}^{5/2} \frac{dT_{tr}}{dt} \quad \left(\eta \equiv 6.5 \times 10^{-29} \frac{L_{\odot}}{M_{\odot}} \right)$$

but $T_{tr} = \left(\frac{L_{\star}}{\eta \mu M_{\star}} \right)^{2/7}$, so

$$\frac{dL_{\star}}{dt} = \frac{7}{2} (\eta \mu)^{2/7} M_{\star}^{2/7} L_{\star}^{5/7} \left[- \frac{L_{\star}}{c_v M_{\star}} \right]$$

$\underbrace{\hspace{10em}}_{\frac{dT_{tr}}{dt} \text{ from specific heat}}$

7.

Or

$$\frac{dL_*}{dt} = -\frac{7}{2} (\eta \mu)^{2/7} M_*^{-5/7} L_*^{5/7+1} c_v^{-1}$$

We can integrate this

$$\int_{L_0}^{L_*} L^{-(5/7+1)} dL = -\frac{7}{2} \frac{(\eta \mu)^{2/7}}{M_*^{5/7} c_v} \int_0^{t_{cool}} dt$$

$$-\frac{7}{5} L^{-5/7} \Big|_{L_0}^{L_*} = -\frac{7}{2} \frac{(\eta \mu)^{2/7}}{M_*^{5/7} c_v} t_{cool}$$

$$\therefore t_{cool} = \frac{2}{5} c_v \frac{M_*^{5/2}}{(\eta \mu)^{2/7}} L_0^{-5/7} \left[\left(\frac{L_*}{L_0} \right)^{-5/7} - \left(\frac{L_0}{L_0} \right)^{-5/7} \right]$$

putting in $\mu = 6.5 \times 10^{-29} \frac{L_0}{M_\odot}$ and $c_v = \frac{1.25 \times 10^8}{\mu_I} \text{ erg/g/K}$

we find

$$t_{cool} \sim 6 \times 10^6 \text{ yr} \left(\frac{A}{12} \right)^{-1} \left(\frac{M}{M_\odot} \right)^{5/7} \left(\frac{\mu}{2} \right)^{-2/7} \left[\left(\frac{L_*}{L_0} \right)^{-5/7} - \left(\frac{L_0}{L_0} \right)^{-5/7} \right]$$

$\mu_I = A$

negligible after long times